



Convective Heat and Mass Transfer Flow of Nanofluid past a Vertical Wavy Wall with Thermal Radiation, Chemical Reaction with Space and Temperature Dependent Heat Sources

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Abstract: Coupled effect of thermal radiation, heat sources and chemical reaction on natural convective heat and mass transfer flow over a vertical wavy surface embedded in a fluid saturated porous medium is considered. The governing equations for flow heat and mass transfer are transformed to a plane geometry case by employing the Runge-Kutta fourth order with shooting technique. The non-dimensional velocity, temperature and nanoconcentration graphs as well as skin friction, rate of heat and mass transfer coefficients are displayed for different values of radiation parameter, heat source parameter, chemical reaction and amplitude of the wavy surface.

Keywords: Nanofluid, Wavy Wall, Thermal Radiation, Chemical reaction, Heat Sources.

1. INTRODUCTION

Magnetic nanofluid has both the liquid and magnetic properties. These fluids have been used in magneto-optical wavelength filters, optical gratings, optical switches etc. The study of Hall effects on nanofluid has applications in MHD energy generators. Study on the problem of Hall currents in nanofluids over a stretching sheet is very limited. Several authors Abo-Eldahab and Abd El-Aziz [1], Ali et al. [3], Abo-Eldahab and Salem [2], Xiaohong Su and Liancun Zheng [15], Mohamed Abd El-Aziz [8]] investigated the effect of Hall currents on convective heat transfer flow of nanofluid past stretching sheet und varied conditions.

In recent years energy and material saving considerations have prompted an expansion of the efforts at producing efficient heat exchanger equipment through augmentation of heat transfer. Hence it is advantageous to go for converging-diverging geometries for improving the design of heat transfer equipment.

Several authors [Rees & Pop[11,12], Ching-Yang Cheng [4], Hossain et al [6], Hady et al [5], Mallikarjuna [7], Vendabai [14], Ramakrishna and Satyanarayana [9], Ranganatha reddy et al[10]] have discussed heat and mass transfer fluid flow past wavy surface at different condition.

Sulochana et al [13] have analysed MHD flow over a permeable stretching /shrinking sheet of a nanofluid in porous medium with suction/injection along with viscous dissipation, considering nanofluid volume fraction on the boundary is submissive to controlled rather than active controlled.

In this paper, we investigate effect of thermal radiation, heat sources, chemical reaction on natural non-darcy convective heat and mass transfer flow over a vertical wavy surface embedded in a fluid saturated porous medium along with viscous dissipation. Here we considered nanofluid volume fraction on the boundary is submissive controlled rather than active controlled.. The vertical wavy wall and the governing equations for flow heat and mass transfer are transformed to a plane geometry case by employing the Runge-Kutta fourth order with shooting technique. The non-dimensional velocity, temperature and concentration graphs as well as skin friction, rate of heat and mass transfer coefficients are displayed for different values of thermal radiation, heat source, chemical reaction parameter and amplitude of the wavy surface.

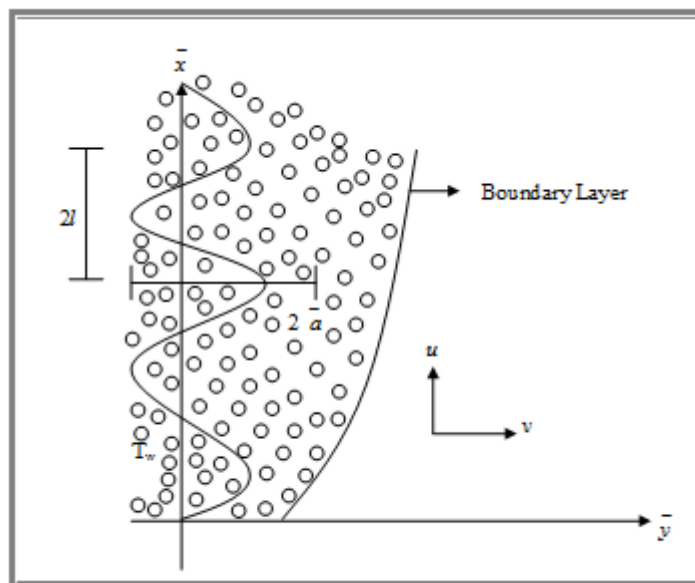


Fig. - 1 : Physical Configuration and Co-ordinate System

2. FORMULATION OF THE PROBLEM

We analyse steady incompressible two-dimensional laminar natural convective heat and mass transfer flow of a nanofluid over a vertical wavy surface embedded in a saturated porous medium. The porous medium is uniform and local thermal equilibrium with the fluid. The Darcy law is used to describe the fluid saturated porous medium. The fluid is assumed to be gray absorbing-emitting radiation but non-scattering medium.

The wavy surface profile is given by

$$y = \bar{\sigma}(\bar{x}) = \bar{a} \sin\left(\frac{\pi \bar{x}}{l}\right) \quad (1)$$

where l is the characteristic length of wavy surface and \bar{a} is the amplitude of the wavy surface. The wavy surface is maintained at constant temperature T_w which are higher than the ambient fluid temperature T_∞ . We consider the natural convection-radiation flow in the presence of heat sources to be governed by the following equations under Boussinesq approximations:

$$\frac{\partial u}{\partial x} + \frac{\partial v}{\partial y} = 0 \quad (2)$$

$$u = -\frac{k}{\mu} \frac{\partial p}{\partial x} + \frac{kg(1-C_\infty)\beta\rho_{f_s}\beta g(T-T_\infty)}{\mu} - \frac{(\rho_p - \rho_{f_s})kg(C-C_\infty)}{\mu} - \frac{k\sigma\mu_e^2 H_0^2}{\mu} u - \left(\frac{b}{k}\right)u^2 \quad (3)$$

$$v = -\frac{k}{\mu} \frac{\partial p}{\partial y} - \left(\frac{b}{k}\right)v^2 \quad (4)$$

$$u \frac{\partial T}{\partial x} + v \frac{\partial T}{\partial y} = \frac{k_f}{\rho_f C_p} \frac{\partial^2 T}{\partial y^2} + \tau \left\{ D_B \left(\frac{\partial T}{\partial x} \frac{\partial C}{\partial x} + \frac{\partial T}{\partial y} \frac{\partial C}{\partial y} \right) + \frac{D_T}{T_\infty} \left[\left(\frac{\partial T}{\partial x} \right)^2 + \left(\frac{\partial T}{\partial y} \right)^2 \right] \right\} - \left(\frac{1}{\rho C_p} \right) \frac{\nabla \cdot q_R}{\partial y} + \frac{q'''}{\rho C_p} \quad (5)$$

$$u \frac{\partial C}{\partial x} + v \frac{\partial C}{\partial y} = D_B \left(\frac{\partial^2 C}{\partial x^2} + \frac{\partial^2 C}{\partial y^2} \right) + \frac{D_T}{T_\infty} \left(\frac{\partial^2 T}{\partial x^2} + \frac{\partial^2 T}{\partial y^2} \right) - kc(C-C_\infty) \quad (6)$$

The relevant boundary conditions are

$$\bar{u} = 0, \bar{v} = 0, T = T_w, D_B \frac{\partial C}{\partial y} + k_f \frac{\partial T}{\partial y} = 0 \text{ at } \bar{y} = \bar{\sigma}(\bar{x}) = \bar{a} \sin\left(\frac{\pi \bar{x}}{l}\right)$$

$$\bar{u} = 0, T \rightarrow T_\infty, C \rightarrow C_\infty \text{ as } \bar{y} \rightarrow \infty \quad (7)$$

Where \bar{u} and \bar{v} are the volume averaged velocity components in the directions of x and y respectively, T, C are temperature, Concentration respectively, ρ is the density of the fluid, μ is the dynamic viscosity of the fluid, k is the permeability of the porous medium, σ is the electrical conductivity, μ_e is the magnetic permeability, H_0 is the strength of the magnetic field, D_B is the Brownian diffusion coefficient and D_T is the thermophoretic diffusion coefficient, $\tau = \frac{(\rho C_p)}{(\rho C_f)}$ is the ratio between the effective heat capacity of the nanoparticle material and heat capacity of the fluid, kc is the coefficient of chemical reaction, β_o are the coefficients of thermal expansion, α is the thermal conductivity, q_r is the radiative heat flux, g is the acceleration due to gravity, q''' is rate of internal heat generation (>0) or absorption (<0) coefficient and Q_H is the strength of the heat source.

Eliminating pressure from equations (3) & (4) we get

$$(1 + M^2) \frac{\partial u}{\partial y} - \frac{\partial v}{\partial x} = \frac{kg(1 - C_\infty)\beta\rho_{f_\infty}\beta g}{\mu} \frac{\partial T}{\partial y} - \frac{(\rho_p - \rho_{f_\infty})kg}{\mu} \frac{\partial C}{\partial y} - \left(\frac{2b}{k}\right) \left(u \frac{\partial u}{\partial y} - v \frac{\partial v}{\partial x}\right) \quad (8)$$

where $M^2 = \frac{k\sigma\mu_e^2 H_0^2}{\mu}$ is the magnetic field parameter.

On using Rosseland approximation the energy equation with non-uniform heat sources reduces to

$$\begin{aligned} \rho C_p (\bar{u} \frac{\partial T}{\partial x} + \bar{v} \frac{\partial T}{\partial y}) &= k_f \left(\frac{\partial^2 T}{\partial x^2} + \frac{\partial^2 T}{\partial y^2} \right) + \frac{16\sigma^* T_\infty^3}{3\beta_R} \left(\frac{\partial^2 T}{\partial x^2} + \frac{\partial^2 T}{\partial y^2} \right) + \\ &+ \frac{kRa}{L^2 \xi} [A^*(T_w - T_\infty)f' + B^*(T - T_\infty)] / \rho C_p + \frac{D_T}{T_\infty} \left(\left(\frac{\partial T}{\partial x} \right)^2 + \left(\frac{\partial T}{\partial y} \right)^2 \right) \end{aligned} \quad (9)$$

On introducing the following non-dimensional variables

$$x = \frac{\bar{x}}{l}, y = \frac{\bar{y}}{l}, a = \frac{\bar{a}}{l}, \sigma = \frac{\bar{\sigma}}{l}, \psi^* = \frac{\bar{\psi}}{l}, \theta = \frac{T - T_\infty}{T_w - T_\infty}, \phi = \frac{C - C_\infty}{C_w - C_\infty}, \bar{u} = \frac{\partial \bar{\psi}}{\partial \bar{y}}, v = -\frac{\partial \bar{\psi}}{\partial \bar{x}} \quad (10)$$

Equations (6),(8) & (9) reduce to

$$\frac{\partial^2 \psi^*}{\partial x^2} + (1 + M^2) \frac{\partial^2 \psi^*}{\partial y^2} = Ra \left(\frac{\partial \theta}{\partial y} - N \frac{\partial \phi}{\partial y} \right) - fs D^{-1} \left(\frac{\partial \psi^*}{\partial y} \frac{\partial^2 \psi^*}{\partial y^2} - \frac{\partial \psi^*}{\partial x} \frac{\partial^2 \psi^*}{\partial x^2} \right) \quad (11)$$

$$\begin{aligned} \left(1 + \frac{4Rd}{3}\right) \left(\frac{\partial^2 \theta}{\partial \bar{x}^2} + \frac{\partial^2 \theta}{\partial \bar{y}^2} \right) + Nb \left(\frac{\partial \theta}{\partial \bar{x}} \frac{\partial \phi^*}{\partial \bar{y}} + \frac{\partial \theta}{\partial \bar{y}} \frac{\partial \phi^*}{\partial \bar{x}} \right) + Nt \left(\left(\frac{\partial \theta}{\partial \bar{x}} \right)^2 + \left(\frac{\partial \theta}{\partial \bar{y}} \right)^2 \right) + \\ + A_{11} f' + B_{11} \theta = \left(\frac{\partial \theta}{\partial x} \frac{\partial \psi^*}{\partial y} - \frac{\partial \theta}{\partial y} \frac{\partial \psi^*}{\partial x} \right) \end{aligned} \quad (12)$$

$$\frac{1}{Le} \left(\frac{\partial^2 \phi}{\partial \bar{x}^2} + \frac{\partial^2 \phi}{\partial \bar{y}^2} \right) + \left(\frac{Nt}{Nb} \right) \left(\frac{\partial^2 \phi}{\partial \bar{x}^2} + \frac{\partial^2 \phi}{\partial \bar{y}^2} \right) - Le\gamma\phi = - \left(\frac{\partial \phi}{\partial \bar{x}} \frac{\partial \psi^*}{\partial \bar{y}} - \frac{\partial \phi}{\partial \bar{y}} \frac{\partial \psi^*}{\partial \bar{x}} \right) \quad (13)$$

$$\text{where } Ra = \frac{Lg\beta(T_w - T_\infty)(1 - C_\infty)}{a\mu}, \quad v = \frac{\mu_\infty}{\rho}, \quad Rd = \frac{4\sigma^* T_\infty^3}{k_f \beta_R}, \quad M^2 = \frac{\sigma\mu_e^2 H_0^2 l^2}{\mu}, \quad Le = \frac{\nu}{D_B},$$

$$\gamma = \frac{kcL^2}{D_B}, \quad N = \frac{(\rho_p - \rho_{f_\infty})(C_w - C_\infty)}{\rho_{f_\infty}(1 - C_\infty)(T_w - T_\infty)}, \quad N_b = \frac{\tau D_B (C_w - C_\infty)}{a}, \quad N_t = \frac{\tau D_T (T_w - T_\infty)}{a T_\infty}, \quad fs = \frac{2b}{\mu l}$$

is the Darcy-Rayleigh number, kinematic viscosity of the fluid, Radiation parameter, magnetic parameter, Lewis number, chemical reaction parameter, buoyancy ratio, Brownian motion parameter, Thermophoresis parameter, the Forchheimer parameter.

The transformed boundary conditions are

$$\psi^* = 0, \theta = 1, Nb \frac{\partial \phi}{\partial y} + Nt \frac{\partial \theta}{\partial y} = 0 \text{ at } y = a \sin(x)$$

$$\frac{\partial \psi^*}{\partial y} \rightarrow 0, \theta \rightarrow 0, \phi \rightarrow \infty \text{ as } y \rightarrow \infty$$
(14)

We can transform the effect of wavy surface from the boundary conditions into the governing equations by using suitable coordinate transformation with boundary layer scaling for the case of free convection. The Cartesian coordinates (x, y) are transformed into the new variables (ξ, η).

We incorporate the effect of effect of wavy surface and the usual boundary layer scaling into the governing equations (11)-(13) for free convection using the transformations and $Ra \rightarrow \infty$ (i.e boundary layer approximation)

$$x = \xi, \bar{\eta} = \frac{y - a \sin(x)}{\xi^{1/2} Ra^{-1/2}}, \psi^* = Ra^{1/2} \psi$$

These transformations are similar to those presented in for instance *Rees and Pop* [12]. We obtain the following boundary layer equations :

$$(1 + a^2 \cos^2 \xi) \frac{\partial \theta}{\partial \eta} \frac{\partial \psi}{\partial \eta} + (1 + a^2 \cos^2 \xi) \frac{\partial^2 \psi}{\partial \eta^2} = Ra \xi^{1/2} \left(\frac{\partial \theta}{\partial \eta} - N \frac{\partial \phi}{\partial \eta} \right) - M^2 \frac{\partial^2 \psi}{\partial \eta^2}$$

$$- fsD^{-1} (1 + a^3 \cos^3(\xi)) \frac{\partial \psi^*}{\partial \eta} \frac{\partial^2 \psi^*}{\partial \eta^2}$$
(15)

$$\xi^{1/2} \left(\frac{\partial \theta}{\partial \xi} \frac{\partial \psi}{\partial \eta} - \frac{\partial \theta}{\partial \eta} \frac{\partial \psi}{\partial \xi} \right) = (1 + a^2 \cos^2(\xi)) \left(1 + \frac{4Rd}{3} \right) \left(\frac{\partial^2 \theta}{\partial \eta^2} \right) +$$

$$+ Nb \frac{\partial \theta}{\partial \eta} \frac{\partial \phi}{\partial \eta} + Nt \left(\frac{\partial \theta}{\partial \eta} \right)^2 + A_{11} f' + B_{11} \theta$$
(16)

$$\xi^{1/2} Le \left(\frac{\partial \phi}{\partial \xi} \frac{\partial \psi}{\partial \eta} - \frac{\partial \phi}{\partial \eta} \frac{\partial \psi}{\partial \xi} \right) = (1 + a^2 \cos^2(\xi)) \left(\frac{\partial^2 \phi}{\partial \eta^2} + \frac{Nt}{Nb} \frac{\partial^2 \phi}{\partial \eta^2} \right) - Le \gamma(\phi)$$
(17)

3. SOULUTION METHODOLOGY

We now introduce the following similarity variables as

$$\eta = \frac{\bar{\eta}}{(1 + a^2 \cos^2(\xi))}, \psi = \xi^{1/2} f(\eta), \theta = \theta(\eta), \phi = \phi(\eta)$$

In equations (15)-(17) we obtain a system of ordinary differential equations as follows:

$$f'' + \theta f' - \frac{M^2}{(1 + a^2 \cos^2(\xi))} f'' = Ra(\theta' - Nr \phi') - fsd^{-1} \left(\frac{1 + a^3 \cos^3(\xi)}{(1 + a^2 \cos^2(\xi))^2} \right) f f''$$
(18)

$$\left(1 + \frac{4Rd}{3} \right) \theta'' + \frac{1}{2} f \theta' + Nb \theta' \phi' + Nt (\theta')^2 + \frac{A_{11} f' + B_{11} \theta}{(1 + a^2 \cos^2(\xi))}$$
(19)

$$\phi'' + \frac{Le}{2} f \phi' + \left(\frac{Nt}{Nb} \right) \theta'' - Le \gamma(1 + a^2 \cos^2(\xi)) \phi$$
(20)

Where prime denotes differentiation with respect to η .

The corresponding boundary conditions are

$$f = 0 \quad \theta = 1 \quad Nb \frac{\partial \phi}{\partial y} + Nt \frac{\partial \theta}{\partial y} = 0 \quad \text{at } \eta = 0 \tag{21}$$

$$f' \rightarrow 0 \quad \theta \rightarrow 0 \quad \phi \rightarrow 0 \quad \text{as } \eta \rightarrow \infty$$

The main results of practical interest in many applications are heat transfer coefficient mass transfer coefficient at the surface. The drag, heat and mass transfer coefficients are expressed in terms of skin friction, Nusselt and Sherwood numbers C_f , Nu_x , Sh_x . Skin friction, (C_f), Nusselt number (Nu_x) and Sherwood number (Sh_x) are defined in terms of Rax and the amplitude ‘a’ as

$$Nu_x = \frac{xq_w}{\alpha_o(T_w - T_\infty)}, \quad Sh_x = \frac{xm_w}{D_b(C_w - C_\infty)} \tag{22}$$

where q_w is the heat flux on the wavy surface and is defined by

$$q_w = -\alpha_o \bar{n} \cdot \nabla T \quad \text{and} \quad \bar{n} = \left(-\frac{a \cos(\xi)}{\sqrt{1+a^2 \cos^2(\xi)}}, \frac{1}{\sqrt{1+a^2 \cos^2(\xi)}} \right)$$

is the unit normal vector to the wavy surface α_o is the effective porous medium thermal conductivity.

Therefore

$$C_f = \frac{f''(0)(1+a^2 \cos^2(\xi))Ra_x^{1/2}}{(1+M^2+a^2 \cos^2(\xi))}, \quad Nu_\xi = -\frac{\theta'(0)Ra_\xi^{1/2}}{\sqrt{1+a^2 \cos^2(\xi)}}, \quad Sh_\xi = -\frac{\phi'(0)Ra_x^{1/2}}{(1+a^2 \cos^2(\xi))} \tag{23}$$

4. CONPARISON

In the absence of magnetic field ($M=0$) and heat sources ($A_{11}=0$ & $B_{11}=0$), temperature ratio ($A=\theta_w=1$) and chemical reaction ($\gamma=0$) the results are in good agreement with *Vendabai et al* [14]

Table1. Comparison

Parameter	Vendabai et al [14]		Present Results		
	$\tau(0)$	$Nu(0)$	$\tau(0)$	$Nu(0)$	
Ra	2	-0.384567	0.552361	-0.384557	0.552357
	4	-0.413678	0.507823	-0.413665	0.507817
	6	-0.537896	0.205912	-0.537888	0.205909
	10	-0.608976	0.196485	-0.608973	0.196465
A1	-0.2	-0.365436	0.632326	-0.365432	0.632320
	-0.4	-0.467885	0.652293	-0.467880	0.652288
	0.2	-0.105678	0.400191	-0.105674	0.400187
	0.4	-0.135679	0.206583	-0.135672	0.206578
B1	-0.2	-0.345678	0.552361	-0.345675	0.552351
	-0.4	-0.554463	0.806242	-0.554453	0.806238
	0.2	0.300211	0.109467	0.300208	0.109461
	0.4	0.167877	0.334523	0.167872	0.334517
Nb	0.1	-0.505678	0.552361	-0.505675	0.552357
	0.3	-0.323453	0.632362	-0.323452	0.632358
	0.4	-0.267898	0.607824	-0.267888	0.607818
	0.5	-0.203885	0.534872	-0.203880	0.534865
Nt	0.1	-0.299288	0.552361	-0.299287	0.552357
	0.2	-0.324728	0.452362	-0.324724	0.452358
	0.3	-0.367809	0.436086	-0.367802	0.436081
	0.4	-0.399432	0.352556	-0.399427	0.352552

a	0.1	-0.337865	0.552361	-0.337845	0.552359
	0.2	-0.321345	0.468091	-0.321340	0.468088
	0.3	-0.310678	0.371734	-0.310675	0.371726
	0.4	-0.301151	0.282947	-0.301147	0.282936
ξ	$\pi/6$	-0.308474	0.552361	-0.308470	0.552354
	$\pi/4$	-0.299728	0.552456	-0.299724	0.552449
	$\pi/3$	-0.298475	0.552567	-0.298465	0.552558
	$\pi/2$	-0.299466	0.552667	-0.299461	0.552657

5. RESULTS AND DISCUSSION

We follow the convention that the non-dimensional concentration is positive/negative according as the actual concentration(C) is greater/lesser than the ambient concentration (C_∞).While the temperature is positive for all variations.

Figs.2a-2c depicts the variation of velocity, temperature and nano-concentration with buoyancy ratio (N). It can be seen from the profiles that when the molecular buoyancy force dominates over the thermal buoyancy force the velocity enhances in the region (0,1) and reduces in the region (1,4) when the buoyancy forces are in the same direction (fig.2a). The temperature enhances with increase in N (fig.2b).Also the actual nanoconcentration enhances in the region (0,1.0) and reduces in the region (1.0,4.0) with increasing N (fig.2c).

Figs.3a-3c show the variation of velocity, temperature and nano-concentration with the influence of radiation parameter (Rd). From fig.3a we find that the velocity reduces in the region (0,0.5) and enhances in the region (0.5,4.0). This means that the thickness of the momentum boundary layer reduces with increasing values of Rd. Fig.3b represent the temperature with Rd. It can be seen from the profiles that an increase in Rd leads to thickening of the thermal boundary layer.The actual nanoconcentration enhances ion the flow region (0,0.5) and reduces in the remaining region(0.5,4.0).(fig.3c).

Figs.4a-4c represent u, θ and ϕ with Eckert number (Ec). From the profiles we notice that higher the dissipation smaller the velocity, larger the temperature and the nano-concentration in the flow region.

Figs.5a-5c and & 6a-6c represent u, θ and ϕ with space dependent heat source parameter (A1) and temperature dependent heat source parameter (B1). From fig.5a & 6a we find that in the region (0, 0.2) the velocity reduces with $A1>0/B1>0$ and enhances with $A1<0 /B1<0$ and in the remaining region (0.2,4.0) a reversed effect is noticed. The temperature enhances with $A1>0/B1<0$ and reduces with $A1<0/B1<0$ in the entire flow region (fig.5b). From fig.5c we find that in the flow regions (0,0.5) and (5,4.0) the nanoconcentration reduces with $A1>0$ and enhances with $A1<0$ while in the flow region (0.5,2.5) the nanoconcentration enhances with $A1>0/B1>0$,reduces with $A1<0/B1<0$ (figs.6b&6c).

The impact of Nb, Nt qualitatively is same. The effect of the Brownian motion parameter (Nb) and thermophoresis parameter (Nt) is shown in Figs. 7a&8a. The velocity decrease with Nb and enhances with Nt in the region (0,1.0),in the region (1,4) the velocity enhances with Nb, reduces with Nt.An increase in Nb leads to an enhancement in temperature in the entire flow region. It is observed that the effect of the thermophoresis parameter is remarkable on the temperature and concentration of the nanoparticles as expected. As the thermophoresis parameter (Nt) increases from 0.1 to 0.2 the temperature decreases and for higher $Nt \geq 0.3$, it enhances in the flow region (0,1.0) significantly throughout the region (7b & 7c). The nanoconcentration enhances with N bands reducers with Nt in the region (0,1.0) and in the flow region(1.0,4.0), ϕ reduces with Nb, reduces with Nt.

Figs.9a-9c show the variation of velocity, temperature and nanoconcentration with chemical reaction parameter (γ). The velocity reduces in the flow region (0,1.5) and enhances in the region (1.5,4.0) in both degenerating/ generation chemical reaction cases.The temperature increases with $0<\gamma>0$ in the entire flow region.An increase in $\gamma>0$ reduces the nanoconcentration in the entire flow region while an increase in $\gamma<0$, reduces θ and enhances ϕ in the region (1.0,4.0).

From 10a-10c represent the effect of Forcheimer parameter (fs) on the non-dimensional velocity, temperature and nanoparticle concentrtration. Fig.10a shows the variation of velocity with fs. In this case the velocity is found to depreciates in the flow region (0,1.0) and enhances in the region(1,4).

From fig.10b & 10c, we find that as the Forcheimer parameter (f_s) increases the temperature increases and nanoconcentration decreases in the flow region. Thus the inertia and boundary effects increases the temperature and depreciates the nanoparticle concentration. This is due to the thickening of the thermal boundary layer and thinning of the solutal boundary layer as a result of increasing values of Forcheimer parameter (f_s).

Figs.11a-11c represent u , θ and ϕ with amplitude of the wavy surface ' a '. From fig.11a & 11b, we find that an increase in amplitude enhances the velocity and temperature in the entire flow region. The nanoconcentration reduces in the region (0,1.0) and enhance in (1.0,4.0) with increase in ' a ' .

Figs.12a-12d represent the variation of u , g , θ and ϕ with Lewis number Le . We find that an increase in Le increases the thickness of linear, angular velocities, thermal and nanoparticle volume fraction.

Figs.13a-13c depict the variation of u , θ and ϕ with stream wise coordinate (ξ). It can be seen from the profiles that an increase in stream wise coordinate decreases the velocity, temperature and increase nanoconcentration in the flow region (0,1.0) and in the remaining region (1,4), they experience a depreciation .

The skin freiction (C_f) at the wall is represented in table.2 for different variatioins. From the tabular values we find that the magnitude of the skin friction enhances with increase in N . Higher the thermal radiation (R_d) / dissipation (E_c) smaller the skin friction at the wall. Higher the strength of the space/temperature dependent heat source larger magnitude of C_f and smaller C_f at the wall in these of heat sink, C_f reduces in both degenerating / generating chemical reaction cases. An increase in Forcheimer parameter (f_s) / Lewis number (Le) leads to a decay in C_f at the wall, Also C_f grows with N_t and decays with N_b at the wall. An increase in amplitude of the wavy surface ' a ' results in a depreciation in the skin friction at the wall. Also an increase in stream wise coordinate (ξ) leads to a depreciation in the magnitude of C_f at the wall.

The rate of heat transfer (Nu) at the wall is displayed in table.2. From the tabular values we find that rate of heat transfer at the wall increases with increase in buoyancy parameter (N). Higher the thermal radiation (R_d) or dissipation smaller the Nusselt number. The rate of heat transfer increases at the wall with space/temperature dependent heat generating source ($A_1 > 0, B_1 > 0$) and reduces in the case of heat absorbing source ($A_1 < 0, B_1 < 0$). Higher Brownian motion/Forchheimer parameter leads to a depreciation in Nu at the wall. It enhances with increase in N_t . The rate of heat transfer at the wall decreases in both degenerating /generating chemical reaction cases. An increase in Le /amplitude (a) reduces the Nusselt number at the wall. The variation of Nu with stream wise coordinate (ξ) shows that Nu increases with increase in enhances with ' ξ ' .

The rate of mass transfer (Sh) at the wall is displayed in table.2. From the tabular values we find that rate of mass transfer at the wall increases with increase in buoyancy ratio (N). From the values we find that Sh decreases with increase in heat generating source and enhances with heat absorbing source. The Sherwood number reduces with increase in R_d / space dependent heat source at the wall. Higher the dissipation (E_c) smaller the Sherwood number at the wall. The rate of mass transfer at the wall reduces with N_b and enhances with N_t / f_s . An increase in Le depreciates the rate of mass transfer at $\eta=0$. The rate of mass transfer at the wall enhances in both degenerating / generating chemical reaction cases. The variation of Sh with amplitude of the wavy surface (a) and stream wise coordinate (ξ) shows that Sh decreases with increase in ' a ' and enhances with ' ξ ' .

6. CONCLUSION

The non-linear equations governing the flow, heat and mass transfer have been solved by Runge-Kutta-shooting technique. The findings of this analysis are

- 1] When the molecular buoyancy force dominates over the thermal buoyancy force the velocity enhances in (0,1) and reduces when the buoyancy forces are in the same direction. The temperature enhances with increase in N . The nanoconcentration enhances in the region (0,4.0) with increasing N_r . An increase in buoyancy ratio (N) enhances the skin friction, Sherwood number and the Nusselt number on $\eta=0$.
- 2] Higher the thermal radiation (R_d) smaller the velocity in the region (0,0.2) and larger in the region (0.2,4.0). An increase in R_d leads to thickening of the thermal boundary layer, thickening of the solutal boundary layer in (0,0.5), thinning of the solutal boundary layer in (0.5,1.0) which results in an enhancement of the temperature and depreciation in the nano-concentration in the flow

region. Higher the thermal radiation smaller the skin friction, the Nusselt and Sherwood numbers on the wall.

- 3] An increase in the strength of the space temperature dependent heat generating source ($A_1 > 0$, $B_1 > 0$) reduces the velocity, enhances temperature and nanoconcentration and while they experience reduction in the flow region with increasing strength of the absorbing heat source ($A_1 < 0$, $B_1 < 0$). Skin friction, Sherwood number and Sherwood number enhances, with $A_1 > 0$, $B_1 > 0$ while they depreciates with $A_1 < 0$, $b_1 < 0$ on $\eta = 0$
- 4] The velocity reduces in the flow region (0,1.5) and enhances in (1.5,4.0) in both degenerating and generating chemical reaction cases. An increase in γ enhances the temperature in the entire flow region. The nanoconcentration reduces in the flow (0,1) in both degenerating / generating chemical reaction cases. In the flow region (1,4) the concentration of the nanoparticles reduces with $\gamma > 0$ and enhances with $\gamma < 0$. The skin friction, Nusselt and Sherwood numbers reduce on $\eta = 0$ in the degenerating generating chemical reaction cases .
- 5] The velocities decreases with thermophoresis parameter (N_b) & enhances Brownian motion parameter (N_t) in (0,1) and in (1,4), u enhances with N_b and reduces with N_t . The temperature and nanoconcentration enhances in the region (0,1) and in the region (1,4) the temperature enhances, the nanoconcentration reduces with N_b with Skin friction, Nu , Sh enhances with N_t and reduces with N_b on the wall $\eta = 0$.
- 6] Higher the dissipative force smaller the velocity, larger temperature and the nanoconcentration in the flow region. Skin friction, Nusselt number Sherwood number reduces with higher values of Ec on $\eta = 0$.
- 7] An increase in f_s reduces the velocity and nanoconcentration in (0,1) and in the region (1,4) velocity enhances, concentration reduces with f_s . The temperature enhances in the entire flow region with f_s . Thus the inertia and boundary effects increases the temperature and depreciates the nanoparticle concentration. This is due to the thickening of the thermal boundary layer and thinning of the solutal boundary layer as a result of increasing values of Forcheimer parameter (f_s). Skin friction, Nusselt number, Sherwood number reduces with higher values of f_s on $\eta = 0$.
- 8] An increased in Lewis number (Le) reduces the velocity, concentration depreciate with Le on $\eta = 0$.
- 9] An increase in amplitude (a) enhances the velocity and temperature. The concentration reduces in the region (0,1) and enhances in (1,4) with increase in ' a '. An increase in amplitude (a) deprecates Nuselt and Sherwood numbers and enhances the skin friction on the wall $\eta = 0$.
- 10] An increase in stream wise coordinate (ξ) decreases the velocity, temperature while enhances the concentration in the flow region (0,1) and depreciates in (1,4). An increase in stream wise coordinate (ξ) reduces skin friction and enhances Nusselt and Sherwood numbers on $\eta = 0$.

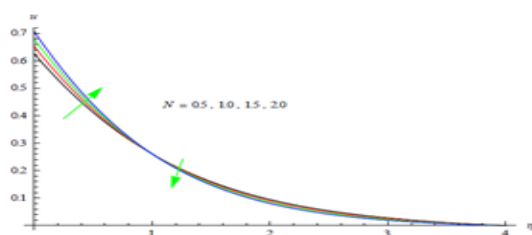


Fig. 2a Variation of velocity(f) with N
 $Rd=0.5, Ec=0.01, Nb=0.3,$
 $Nt=0.3, A_1=0.5, B_1=0.5, a=0.1, \xi=\pi/4$

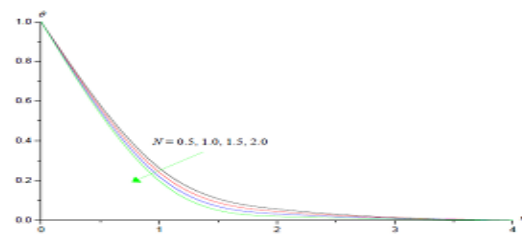


Fig. 2b Variation of temperature(θ) with N
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 $Nt=0.3, A_1=0.5, B_1=0.5, a=0.1, \xi=\pi/4$

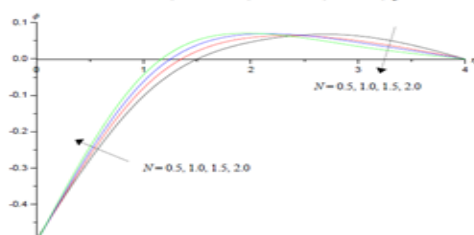


Fig. 2c Variation of nanoconcentration(ϕ) with N
 $Rd=0.5, Ec=0.01, Nb=0.3,$
 $Nt=0.3, A_1=0.5, B_1=0.5, a=0.1, \xi=\pi/4$

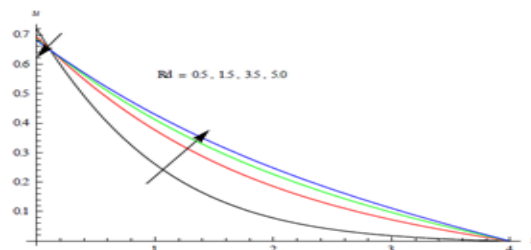


Fig. 3a Variation of velocity(f) with Rd
 $N=0.5, Ec=0.01, Nb=0.3,$
 $Nt=0.3, A_1=0.5, B_1=0.5, a=0.1, \xi=\pi/4$

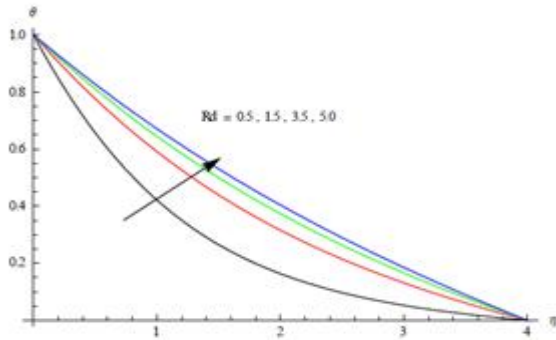


Fig.3b Variation of temperature(θ) with Rd
 $N=0.5, Ec=0.01, Nb=0.3,$
 $Nt=0.3, A1=0.5, B1=0.5, a=0.1, \xi=\pi/4$

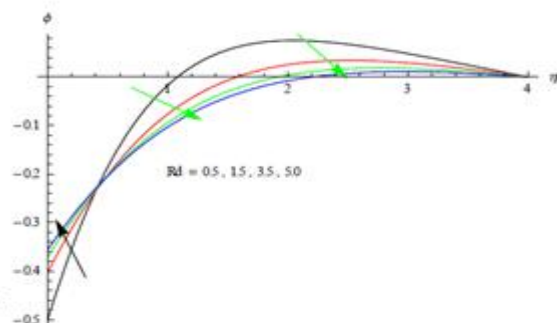


Fig.3c Variation of nanoconcentration(ϕ) with Rd
 $N=0.5, Ec=0.01, Nb=0.3,$
 $Nt=0.3, A1=0.5, B1=0.5, a=0.1, \xi=\pi/4$

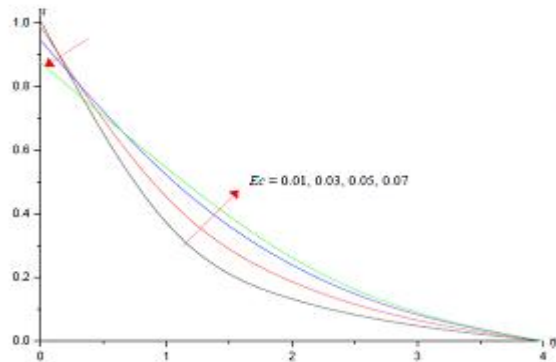


Fig.4a Variation of velocity(f) with Ec
 $N=0.5, Rd=0.5, Nb=0.3,$
 $Nt=0.3, A1=0.5, B1=0.5, a=0.1, \xi=\pi/4$

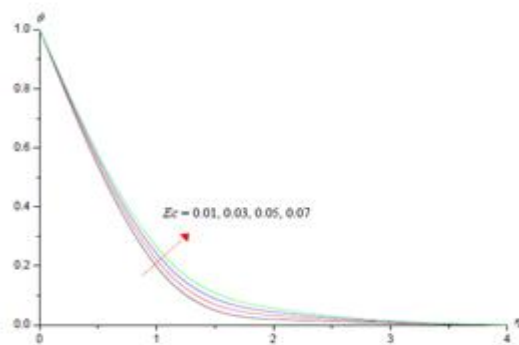


Fig.4b Variation of temperature(θ) with Ec
 $N=0.5, Rd=0.5, Nb=0.3, Nt=0.3,$
 $A1=0.5, B1=0.5, a=0.1, \xi=\pi/4, Le=2$

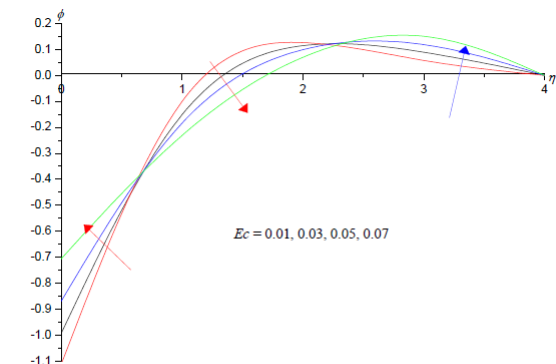


Fig.4c Variation of nanoconcentration(ϕ) with Ec
 $N=0.5, Rd=0.5, Nb=0.3, Nt=0.3,$
 $A1=0.5, B1=0.5, a=0.1, \xi=\pi/4, Le=2$

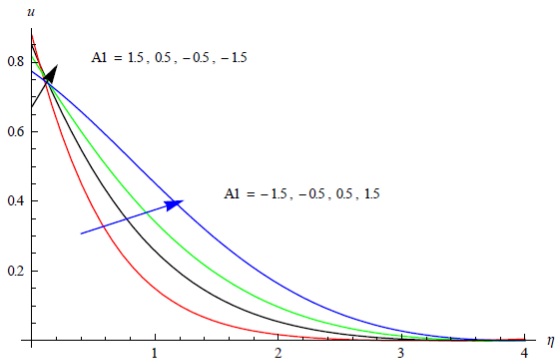


Fig.5a Variation of velocity(f) with A1
 $N=0.5, Rd=0.5, Ec=0.01, Nb=0.3,$
 $Nt=0.3, B1=0.5, a=0.1, \xi=\pi/4$

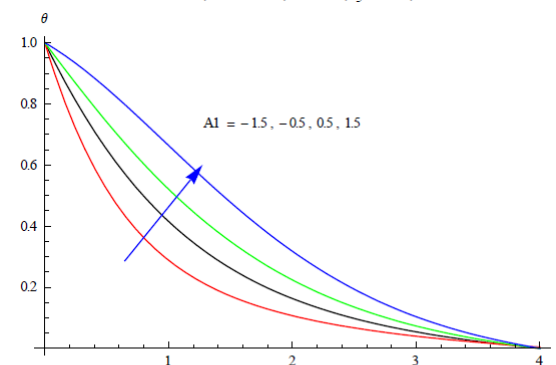


Fig.5b Variation of temperature(θ) with A1
 $N=0.5, Rd=0.5, Ec=0.01, Nb=0.3,$
 $Nt=0.3, B1=0.5, a=0.1, \xi=\pi/4$

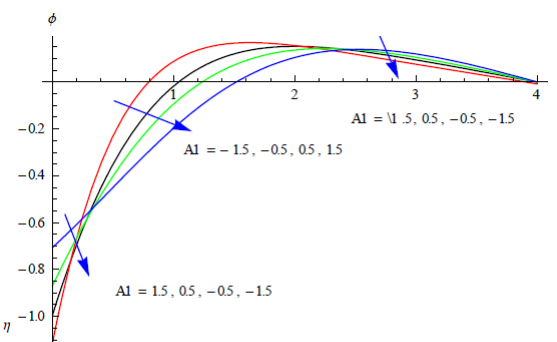


Fig.5c Variation of nanoconcentration(ϕ) with A1
 $N=0.5, Rd=0.5, Ec=0.01, Nb=0.3,$
 $Nt=0.3, B1=0.5, a=0.1, \xi=\pi/4$

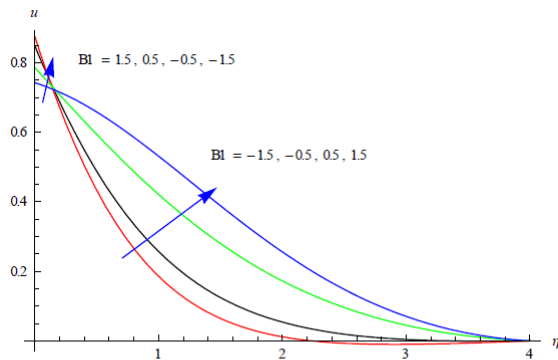


Fig.6a Variation of velocity(f') with $B1$
 $N=0.5, Rd=0.5, Ec=0.01, Nb=0.3,$
 $Nt=0.3, A1=0.5, a=0.1, \xi=\pi/4$

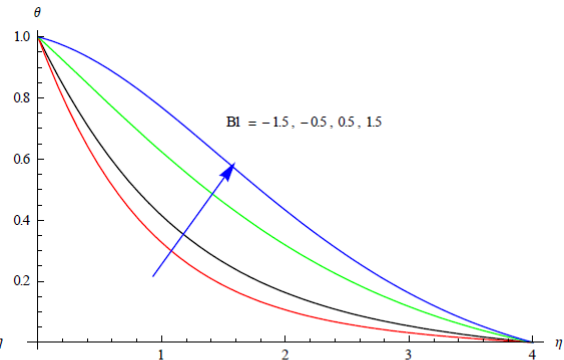


Fig.6b Variation of temperature(θ) with $B1$
 $N=0.5, Rd=0.5, Ec=0.01, Nb=0.3,$
 $Nt=0.3, A1=0.5, a=0.1, \xi=\pi/4$

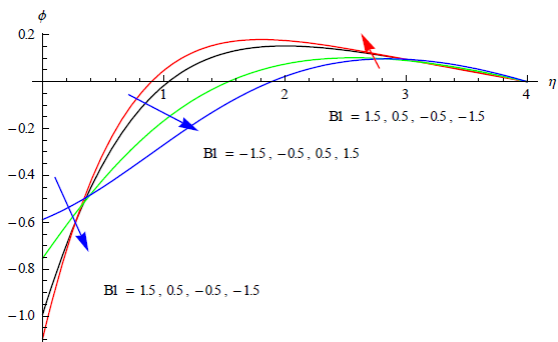


Fig.6c Variation of nanoconcentration(ϕ) with $B1$
 $N=0.5, Rd=0.5, Ec=0.01, Nb=0.3,$
 $Nt=0.3, A1=0.5, a=0.1, \xi=\pi/4$

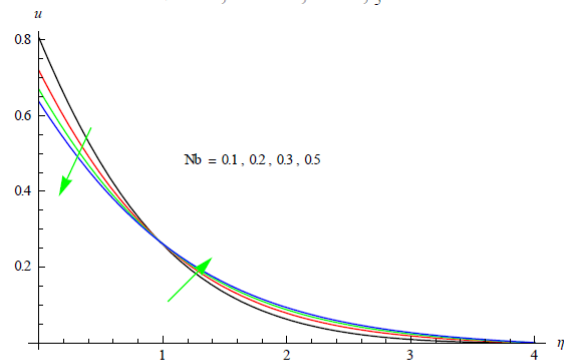


Fig.7a Variation of velocity(f') with Nb
 $N=0.5, Rd=0.5, Ec=0.01,$
 $Nt=0.3, A1=0.5, B1=0.5, a=0.1, \xi=\pi/4$

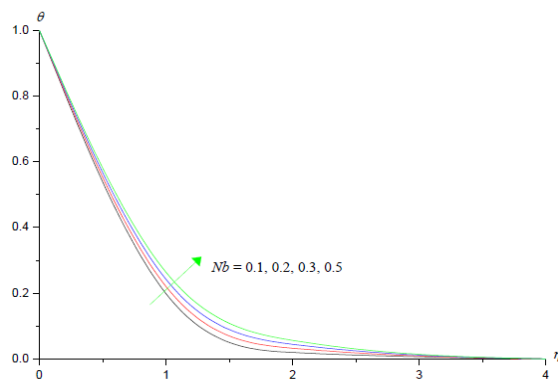


Fig.7b Variation of temperature(θ) with Nb
 $N=0.5, Rd=0.5, Ec=0.01, Nt=0.3, A1=0.5,$
 $B1=0.5, a=0.1, \xi=\pi/4$

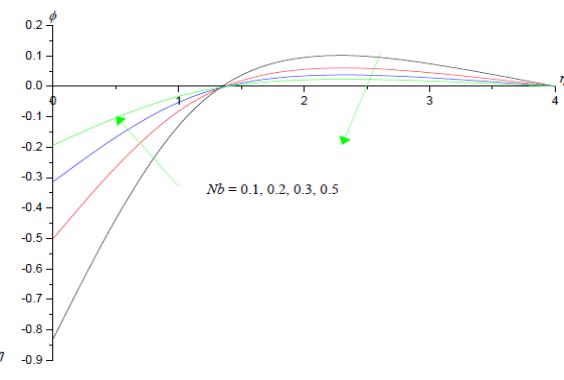


Fig.7c Variation of nanoconcentration(ϕ) with Nb
 $N=0.5, Rd=0.5, Ec=0.01, Nt=0.3,$
 $A1=0.5, B1=0.5, a=0.1, \xi=\pi/4$

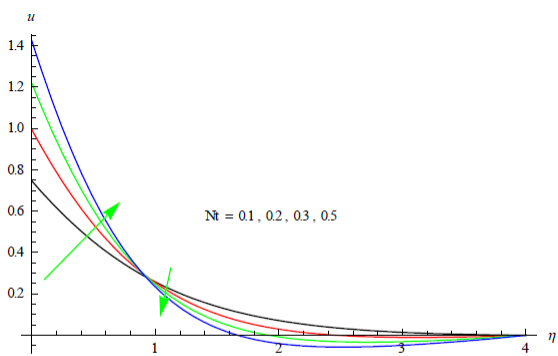


Fig.8a Variation of velocity(f') with Nt
 $N=0.5, Rd=0.5, Ec=0.01, Nb=0.3,$
 $A1=0.5, B1=0.5, a=0.1, \xi=\pi/4$

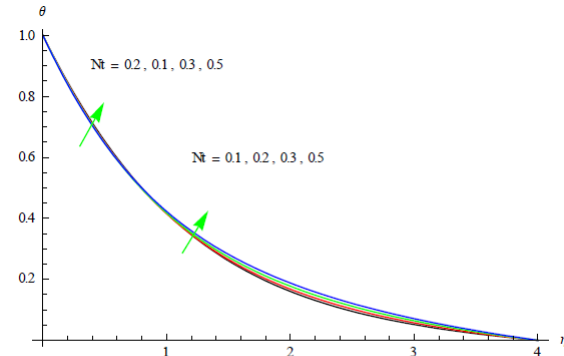


Fig.8b Variation of temperature(θ) with Nt
 $N=0.5, Rd=0.5, Ec=0.01, Nb=0.3,$
 $A1=0.5, B1=0.5, a=0.1, \xi=\pi/4$

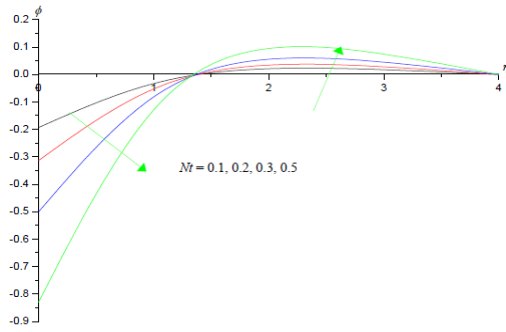


Fig. 8c Variation of nanoconcentration(ϕ) with Nt
 $N=0.5, Rd=0.5, Ec=0.01, Nb=0.3,$
 $A1=0.5, B1=0.5, a=0.1, \xi=\pi/4$

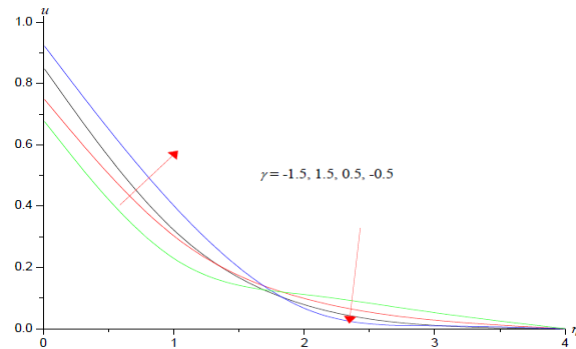


Fig. 9a Variation of velocity(f') with γ
 $N=0.5, Rd=0.5, Ec=0.01, Nb=0.3,$
 $Nt=0.3, A1=0.5, B1=0.5, a=0.1, \xi=\pi/4$

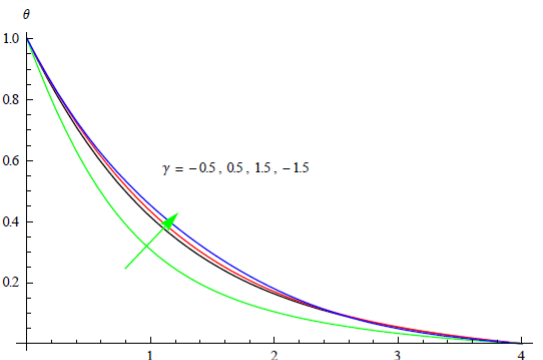


Fig. 9b Variation of temperature(θ) with γ
 $N=0.5, Rd=0.5, Ec=0.01, Nb=0.3,$
 $Nt=0.3, A1=0.5, B1=0.5, a=0.1, \xi=\pi/4$

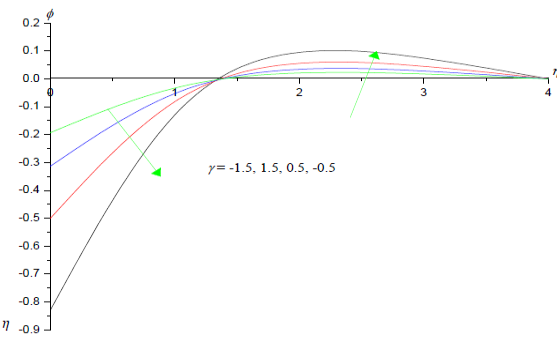


Fig. 9c Variation of nanoconcentration(ϕ) with γ
 $N=0.5, Rd=0.5, Ec=0.01, Nb=0.3,$
 $Nt=0.3, A1=0.5, B1=0.5, a=0.1, \xi=\pi/4$

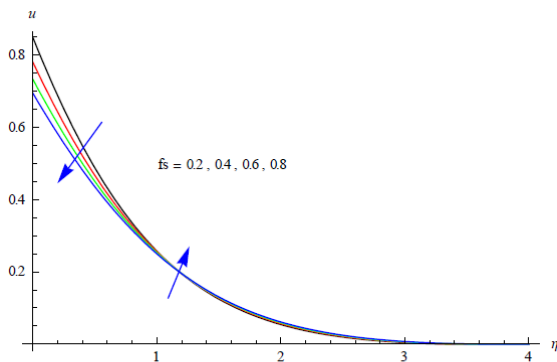


Fig. 10a Variation of velocity(f') with f_s
 $N=0.5, Rd=0.5, Ec=0.01, Nb=0.3,$
 $Nt=0.3, A1=0.5, B1=0.5, a=0.1, \xi=\pi/4$

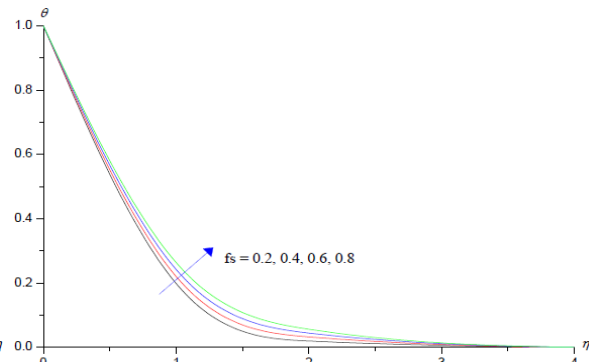


Fig. 10b Variation of temperature(θ) with f_s
 $N=0.5, Rd=0.5, Ec=0.01, Nb=0.3,$
 $Nt=0.3, A1=0.5, B1=0.5, a=0.1, \xi=\pi/4$

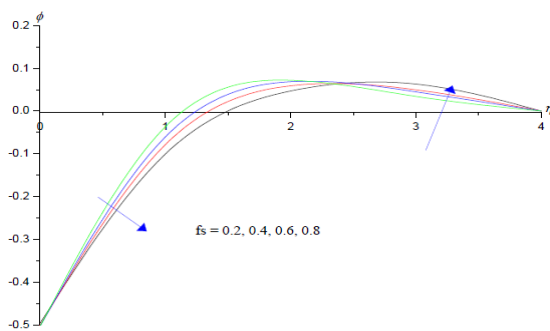


Fig. 10c Variation of nanoconcentration(ϕ) with f_s
 $N=0.5, Rd=0.5, Ec=0.01, Nb=0.3,$
 $Nt=0.3, A1=0.5, B1=0.5, a=0.1, \xi=\pi/4$

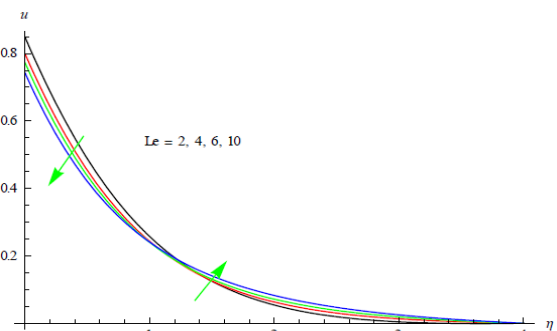


Fig. 11a Variation of velocity(f') with Le
 $N=0.5, Rd=0.5, Ec=0.01, Nb=0.3,$
 $Nt=0.3, A1=0.5, B1=0.5, a=0.1, \xi=\pi/4$

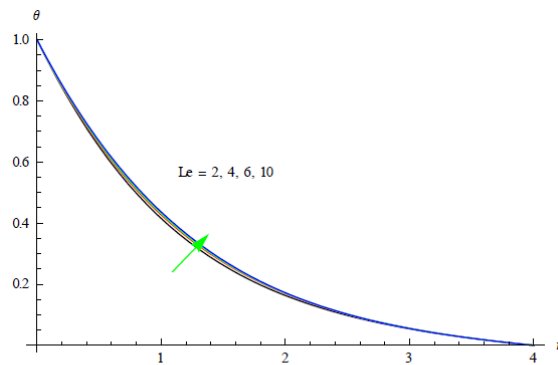


Fig. 11b Variation of temperature(θ) with Le
 $N=0.5, Rd=0.5, Ec=0.01, Nb=0.3,$
 $Nt=0.3, A1=0.5, B1=0.5, a=0.1, \xi=\pi/4$

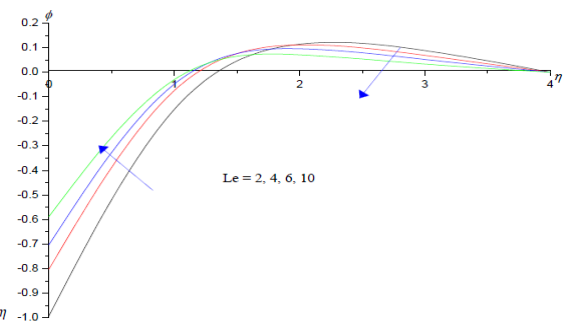


Fig. 11c Variation of nanoconcentration(ϕ) with Le
 $N=0.5, Rd=0.5, Ec=0.01, Nb=0.3,$
 $Nt=0.3, A1=0.5, B1=0.5, a=0.1, \xi=\pi/4$

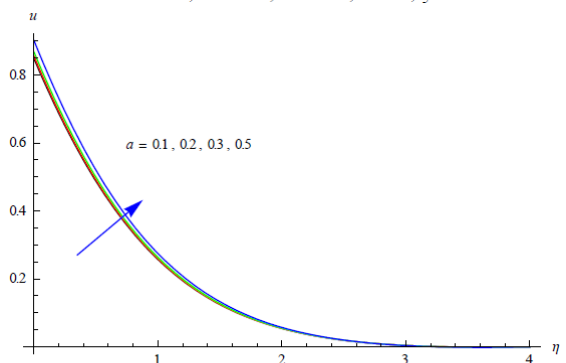


Fig. 12a Variation of velocity(f) with a
 $N=0.5, Rd=0.5, Ec=0.01, Nb=0.3,$
 $Nt=0.3, A1=0.5, B1=0.5, a=0.1, \xi=\pi/4$

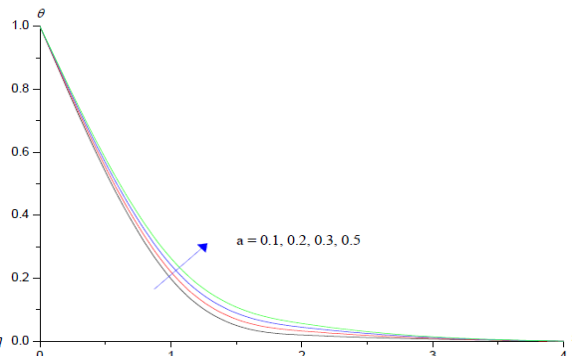


Fig. 12b Variation of temperature(θ) with a
 $N=0.5, Rd=0.5, Ec=0.01, Nb=0.3,$
 $Nt=0.3, A1=0.5, B1=0.5, a=0.1, \xi=\pi/4$

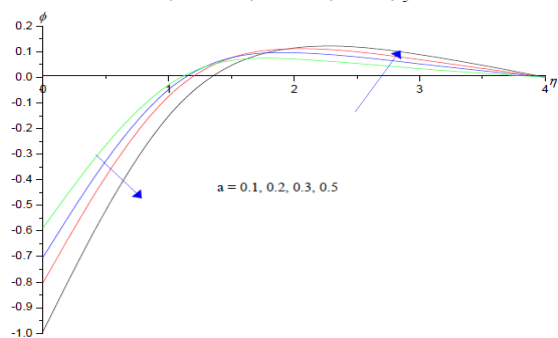


Fig. 12c Variation of nanoconcentration(ϕ) with a
 $N=0.5, Rd=0.5, Ec=0.01, Nb=0.3,$
 $Nt=0.3, A1=0.5, B1=0.5, a=0.1, \xi=\pi/4$

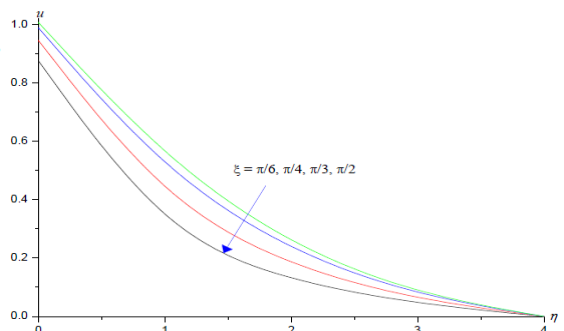


Fig. 13a Variation of velocity(f) with ξ
 $N=0.5, Rd=0.5, Ec=0.01, Nb=0.3,$
 $Nt=0.3, A1=0.5, B1=0.5, a=0.1, \xi=\pi/4$

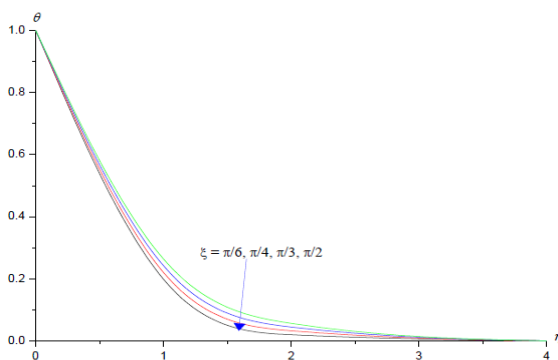


Fig. 13b Variation of temperature(θ) with ξ
 $N=0.5, Rd=0.5, Ec=0.01, Nb=0.3,$
 $Nt=0.3, A1=0.5, B1=0.5, a=0.1$

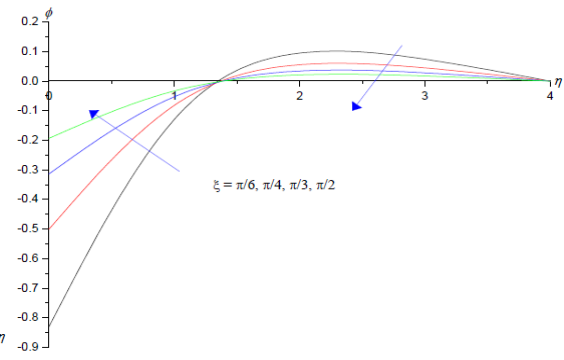


Fig. 13c Variation of nanoconcentration(ϕ) with ξ
 $N=0.5, Rd=0.5, Ec=0.01, Nb=0.3,$
 $Nt=0.3, A1=0.5, B1=0.5, a=0.1$

Table2. Skin Friction (τ_x), Nusslet number (Nu) and Sherwood Number (Sh) at $\eta = 0$

Parameter	$\tau_x(0)$	$Nu(0)$	$Sh(0)$	Parameter	$\tau_x(0)$	$Nu(0)$	$Sh(0)$		
N	0.5	-0.494305	-0.494305	-0.792374	B1	0.5	-0.861994	0.832395	-1.66479
	1	-0.537661	-0.537661	-0.797519		1.5	-1.1163	1.08385	-2.1677
	1.5	-0.581011	-0.581011	-0.80254		-0.5	-0.39144	0.373203	-0.746407
	2	-0.624345	-0.624345	-0.80744		-1.5	-0.110211	0.104128	-0.208256
Rd	0.5	-0.646003	0.809846	-0.809846	γ	0.5	-0.861994	0.832395	-1.66479
	1.5	-0.392428	0.48927	-0.48927		1.5	-0.831937	0.787047	-1.57409
	3.5	-0.324659	0.404067	-0.404067		-0.5	-1.02049	1.08202	-2.16404
	5	-0.292584	0.363828	-0.363828		-1.5	-0.825406	0.780076	-1.56015
Ec	0.01	-0.861994	0.832395	-1.66479	fs	0.2	-0.861994	0.832395	-1.66479
	0.03	-0.861241	0.83166	-1.66332		0.4	-0.729958	0.818463	-1.63693
	0.05	-0.860524	0.830959	-1.66192		0.6	-0.64856	0.8081	-1.6162
	0.07	-0.859806	0.830257	-1.66051		0.8	-0.587976	0.799267	-1.59853
Nb	0.1	-0.790142	0.825171	-1.37529	a	0.1	-0.861994	0.832395	-1.66479
	0.2	-0.646003	0.809846	-0.809846		0.2	-0.869177	0.829568	-1.65914
	0.3	-0.564756	0.800672	-0.50042		0.3	-0.87935	0.825361	-1.65072
	0.5	-0.512648	0.794566	-0.305602		0.5	-0.907049	0.81296	-1.62592
Nt	0.2	-0.694279	0.820572	-0.984686	ξ	$\pi/6$	-0.86308	0.831978	-1.66396
	0.4	-1.10769	0.847242	-2.71118		$\pi/4$	-0.861994	0.832395	-1.66479
	0.6	-1.50101	0.866015	-4.50328		$\pi/3$	-0.860893	0.832813	-1.66563
	0.8	-1.87472	0.879401	-6.33169		$\pi/2$	-0.859773	0.833231	-1.66646
A1	0.5	-0.861994	0.832395	-1.66479					
	1.5	-1.37235	1.34164	-2.64923					
	-0.5	-0.566758	0.543703	-1.08741					
	-1.5	-0.247629	0.235471	-0.470942					

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Citation: *Dr. M. Nagasasikala., Convective Heat and Mass Transfer Flow of Nanofluid past a Vertical Wavy Wall with Thermal Radiation, Chemical Reaction with Space and Temperature Dependent Heat Sources, International Journal of Scientific and Innovative Mathematical Research (IJSIMR), vol. 8, no. 9, pp. 22-35, 2020. Available : DOI: <https://doi.org/10.20431/2347-3142.0809004>*

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