



The Molecular Diffusion of Heat and Mass from Two Spheres

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Abstract: *An analytical solution to the heat conduction equation applied to two spheres was obtained. Bispherical coordinates system were used. The derived solution applies to both fluid and solid spheres. The solution revealed a decrease in the heat conduction between the two spheres when they approach one another. The solution compared well with the available theoretical models*

Keywords: *Heat Conduction, Two Spheres, Heat Transfer, Temperature, Analytical Solution*

1. INTRODUCTION

Heat and mass transfer from spheres immersed in fluids occur in many engineering industries, technologies and scientific applications. Among these are: drying, adsorption, extraction, fixed and fluidized beds, cloud physics, aerosol physics, combustion of fuel droplets and cooling of spherical uranium fuel elements in certain types of nuclear reactors.

Although an individual spherical particle or bubble is rare to be found in practical systems, yet most of the information of the scientific literature deals with studies of the individual sphere or bubble. The heat

Conduction from a single sphere to the surrounding mass of fluid is given as follows (Leal 1992)

$$(Nu)_s = 2 \quad (1)$$

The above equation applies also to the case of mass transfer simply by replacing Nu by the Sherwood number Sh. Morrison and Reed (1974) derived a solution for the conductive heat transfer from two touching spheres. They found that the Nusselt number of a sphere was reduced by the presence of a touching sphere according to the following

$$Nu = 2 \ln 2 \quad (2)$$

The problem of two, three and multi-sphere or multi-bubble system are now being addressed by a number of research workers e.g. Ruzicka (2000) and Ramachandran et al. (1989) who solved the forced convective heat transfer to a linear array of three spheres by using the finite element method. In general, an increased rate of convective heat or mass transfer with sphere separation was obtained.

The molecular heat conduction and mass diffusion of the two spheres are governed by the following phenomenon: the diffusion of heat or mass from one sphere moderates the gradient of temperature and concentration around the other sphere and thus Nu or Sh are decreased with the separation distance between spheres. This phenomenon is encountered in low Reynolds number flow (e.g. Chen and Pfeiffer (1970), Aminzadeh et al. (1974) and Ramachandran and Kleinstreuer (1985).

The following conclusions were deduced from the literature review:

- A confusion exists as to the increase or decrease of heat transfer with sphere spacing. (see for example Mulholland et al. (1988).
- The majority of the theoretical research was confined to the numerical solution of the conservation equations, except for the work of Kendoush (2007).
- Two-sphere interactions were not treated experimentally, except for the works of Wang and Liu

(1992) of the two solid spheres, Kok (1989), and Sanada (2005) of the two bubbles.

The ideal case of the present paper serves the purpose of modeling practical situations of conductive flow of heat in a system of two spheres. The two-sphere analysis may be considered as the lowest order effect of thermal interaction of multi-spheres in a swarm. The present presentation may be considered as a complementary to that of Kendoush (2007).

2. THEORETICAL ANALYSIS

Assume two spheres of equal radii a separated by a certain distance L as shown in Fig. 1. Yovanovich (1978) developed a method of dealing with heat conduction in complicated geometries. The method depends on choosing an orthogonal curvilinear coordinate system which is most appropriate for the problem at hand, and solving Laplace's equation in that coordinate system. Accordingly, the bispherical coordinate system was chosen here for the present problem. The heat conduction equation, in bispherical coordinates, is given as follows

$$\frac{(\cosh\eta - \cos\theta)^3}{c^2 \sin\theta} \left[\sin\theta \frac{\partial}{\partial \eta} \left(\frac{1}{\cosh\eta - \cos\theta} \frac{\partial T}{\partial \eta} \right) + \frac{\partial}{\partial \theta} \left(\frac{\sin\theta}{\cosh\eta - \cos\theta} \frac{\partial T}{\partial \theta} \right) \right] = 0 \quad (3)$$

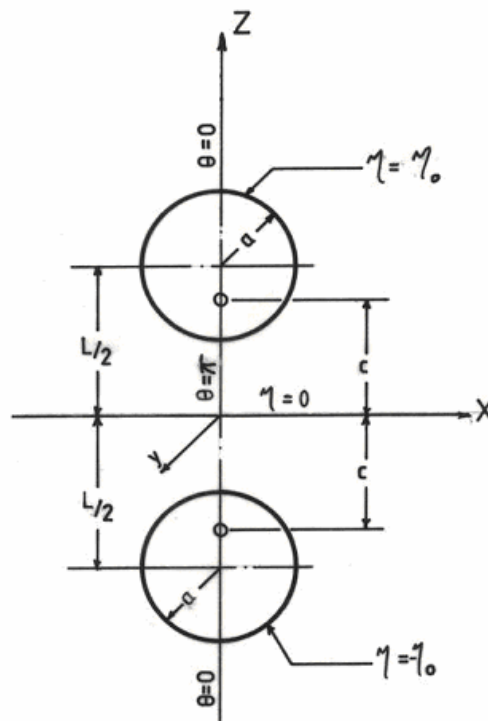


Fig1. The two spheres in bispherical coordinates.

It will be supposed that heat is generated within the two spheres at a rate such that their surface are maintained at a uniform temperature T_a . The boundary conditions, which specify the temperature and heat flow of the system become the following

$$T = T_a \text{ at } \eta = \eta_o \quad (4)$$

$$\frac{\partial T}{\partial \eta} = 0 \text{ at } \eta = 0 \quad (5)$$

$$\frac{\partial T}{\partial \theta} = 0 \text{ at } \theta = 0 - \pi \quad (6)$$

Laplace's Eq. (3) is solved by separation of variables. A solution of sufficient generality for our purpose is the following

$$T = (\cosh\eta - \cos\theta)^{1/2} [B e^{(n+0.5)\eta} + D e^{-(n+0.5)\eta}] P_n(\cos\theta) \quad (7)$$

where B and D are arbitrary constants and $p_n(\cos\theta)$ is the associated Legendre function of the first kind. Applying the boundary conditions makes the complete solution as follows

$$T = T_a [2(\cosh \eta - \cos \theta)]^{0.5} \sum_{n=0}^{\infty} e^{-(n+0.5)\eta_o} \frac{\cosh(n+0.5)\eta}{\cosh(n+0.5)\eta_o} P_n(\cos \theta) \quad (8)$$

Assume that the mass of the stationary fluid round the two spheres has a constant thermal conductivity k .

The heat flux from the spheres surfaces S_o is given by the following

$$q = \iint_{S_o} k g^{1/2} \left(\frac{\partial T}{\partial \eta} \right) n_o dS_n \quad (9)$$

where

$$g_1 = c^2 / (\cosh n - \cos \theta)^2$$

and

$$dS_n = (g_2 g_3)^{1/2} d\theta = \frac{c^2 \sin \theta}{(\cosh n_o - \cos \theta)^2} d\theta \quad (10)$$

The metric coefficients of the coordinates are the following

$$g_1 = g_2 \quad (11)$$

And

$$g_3 = \frac{c^2 \sin^2 \theta}{(\cosh n - \cos \theta)^2} \quad (12)$$

The following dimensions are defined as follows: $a = c / \sinh n_o$ and $L = 2c / \tanh n_o$. The temperature gradient is obtained from Eq. (8) and substituted into Eq. (9) to get the following equation

$$q = 4\pi a k T_a \sum_{n=0}^{\infty} \frac{(-1)^n \sinh n_o}{\sinh(n+1)n_o} \quad (13)$$

Equating the above equation to $q = h(4\pi a^2)T_a$, gives the following equation

$$Nu = 2 \sum_{n=0}^{\infty} \frac{(-1)^n \sinh n_o}{\sinh(n+1)n_o} \quad (14)$$

Or the following

$$Nu = \sum_{n=0}^{\infty} \frac{(-1)^n \sinh \cosh^{-1} \left(\frac{L}{2a} \right)}{\sinh(n+1) \cosh^{-1} \left(\frac{L}{2a} \right)} \quad (15)$$

Equation (15) reduces to the single sphere solution (that is, $Nu=2$ of Eq.(1) upon increasing the value of η_o or L . This increase tends to make the two spheres an infinite distance apart. On the other hand, when $\eta_o=0.2$ (the case of almost touching spheres), we obtain from Eq. (15) the solution of Morrison and Reed [2] (that is, $Nu=2 \ln 2$ of Eq. (2)). Dividing Eq. (15) by the single sphere, we get the following equation

$$Nu / (Nu)_s = \sum_{n=0}^{\infty} (-1)^n \sinh \eta_o / \sinh(n+1)\eta_o \quad (16)$$

3. DISCUSSION

The above analyses of the two spheres apply to both fluid and solid spheres. Figure2 shows that when the centers of the spheres are only 4 radii apart the Nusselt number is approximately 1.6 and when the centers of the spheres are 100 radii apart the Nusselt number is approximately 1% less than the limiting value of 2. It can be deduced from Fig.2 that as the number of spheres is increased the Nusselt number becomes progressively smaller until a condition is reached in which a particular sphere is surrounded by multiple spheres of the same surface temperature. This means that a point will be reached where the Nusselt number approaches zero. Figure 2 shows a

comparison between the present solution and that of Morrison and Reed (1974). Both solutions agree on the decrease of Nu number with separation but the present solution has the following merits

- (i) It is in a closed form. This would enable researchers to utilize the present solution together with the earlier solution of the author (Kendoush (2007)) to obtain the complete solution of heat transfer to two bubbles vertically above one another, as follows

$$Nu = 2 \sum_{n=0}^{\infty} \frac{(-1)^n \sinh \cosh^{-1} \left(\frac{L}{2a} \right)}{\sinh(n+1) \cosh^{-1} \left(\frac{L}{2a} \right)} + \frac{1}{2} \left[\frac{2}{\sqrt{\pi}} \sqrt{H} ([Pe]_A)^{1/2} + \frac{2}{\sqrt{\pi}} \sqrt{G} ([Pe]_B)^{1/2} \right] \quad (17)$$

where

$$H = [1 + (a/L)^3]$$

and

$$G = \left[1 + \left(\frac{a}{L} \right)^6 \right]$$

Subscripts A and B refers to top and bottom bubbles respectively.

- (ii) The present results are closer to those of Russell (1911) than Morrison and Reed (1974). Russell (1911) calculated electric charges on two spherical condensers whose surfaces were at the same potential. His results were presented in approximate formulas and tables.

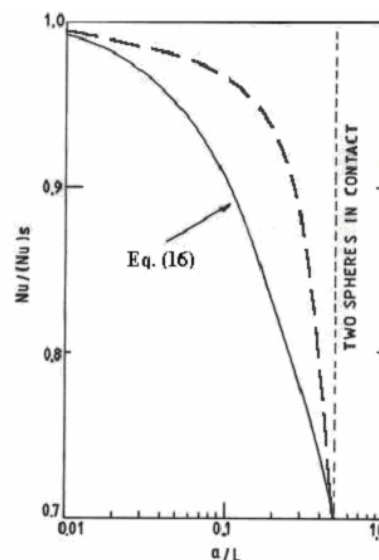


Fig2. Comparison between the present solution (-----) and that of Morrison and Reed (1974) (-----)

4. CONCLUSION

The present paper revealed and proved analytically the following phenomenon:

“A decrease in the heat conduction and mass diffusion between two solid or fluid spheres occurs when they approach one another.” This is contrary to the case of heat convection Kendoush (2007).

In general, the present analytical results compared well with the available theoretical models.

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NOMENCLATURE

a	=	radius of the sphere
h	=	convective heat transfer coefficient
k	=	thermal conductivity of the fluid
L	=	the distance between centers of the two spheres
Nu	=	Nusselt number ($2ah/k$)
q	=	the heat flux on the sphere surface
Re	=	Reynolds number ($2a\rho U/\mu$)
T	=	temperature
T_a	=	absolute temperature of the sphere surface
η, θ	=	bispherical coordinates

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